

## PECULIARITIES OF TIME-DEPENDENT-DIELECTRIC BREAKDOWN CHARACTERISTICS OF PURE AND DOPED TA<sub>2</sub>O<sub>5</sub> STACKS

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**Abstract.** *The effect of both the process-induced defects and the dopant on the time-dependent-dielectric breakdown in Ta<sub>2</sub>O<sub>5</sub> stacks is discussed. The breakdown degradation is analyzed in terms of specific properties of high-*k* stacks which make their dielectric breakdown mechanism completely different from that of classical SiO<sub>2</sub>. The relative impact of a number of factors constituting the reliability issues in Ta<sub>2</sub>O<sub>5</sub>-based capacitors (trapping in pre-existing traps, stress-induced new traps generation, the presence of interface layer at Si and the role of the dopant) is clarified.*

**Keywords:** *high-*k* stack; doped Ta<sub>2</sub>O<sub>5</sub>; Weibull distribution; hard breakdown*

### 1. INTRODUCTION

A considerable progress has been made in understanding the electrical and material properties of a number of high-*k* dielectrics, and the main purpose for their implementation in nanoelectronic devices (low leakage current for equivalent oxide thickness below about 1 nm) has been achieved. In order to fulfil also the criterion of dielectric reliability assurance, a deeper physical understanding of the dielectric breakdown (DB) mechanisms in these films is required. A special focus at present is on doped and multicomponent high-*k* dielectrics as a new step in high-*k* alternatives engineering. A little is known on the processes which are forerunners of the BD events in high-*k* stacks, on the effect of dopant in constituting the pre-existing traps and their impact on the BDs. The specific feature of high-*k* dielectrics, charge trapping, fundamentally defines the reliability of high-*k* stacks and dictates the allowed operating voltage for 10 years lifetime requirement.

The dielectric reliability is typically denoted by charge  $Q_{bd}$  or time  $t_{bd}$  to breakdown, and an important aspect of this reliability is the voltage dependence of these entities. It is generally accepted that MOS devices are required to work constantly for about 10 years. Increased voltage and temperature are commonly used to accelerate the breakdown pro-

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cess, since waiting for 10 years to measure failure would not be possible. One approach for understanding the high- $k$  degradation mechanisms is using indirect methods such as those based on data for stress-induced leakage current (SILC) [1]-[5]. Essential disadvantage of this approach is the fact that SILC appears long before the breakdown events. Whatever the specific approach is, however, a correct way for obtaining adequate picture of the voltage dependence of the dielectric BD is to collect data with great experimental statistics using  $Q_{bd}$  and  $t_{bd}$  as variable entities. It was suggested that the degradation mechanism in high- $k$  dielectrics is dominated by an "extrinsic" mechanism in contrast to so called "intrinsic" BD which is not related to process-induced or other extraneous defects. The intrinsic reliability, typically assigned to  $\text{SiO}_2$ , is a material property and it is sufficiently well understood. Having in mind the inherent property of high- $k$  insulators to be a trap-rich materials the behaviour of the early (extrinsic) failures is very important for these dielectrics. Due to the strong trapping phenomena, however, the relation between the BD and SILC as a forerunner of BD known for  $\text{SiO}_2$ , could be no longer valid in high- $k$  films and BD can occur without occurrence of any significant SILC and vice versa – a significant SILC can be created before the high- $k$  dielectric is broken down [6]. So the pre-existing traps were identified as responsible for the domination of extrinsic BD mechanism in these films. Since the parameters of each high- $k$  dielectric are a function of both its chemical composition and the fabrication process [3], [5], [7]-[11], Time-Dependent-Dielectric Breakdown (TDDB) characteristics are difficult to model from fundamental dielectric properties point of view. This circumstance hampers the interpretation of the degradation mechanism of a certain high- $k$  film by analogy to another one. This imperatively means that the BD processes have to be studied individually for each high- $k$  film. It is expected that from the acquisition of many experimental data some fundamental pattern could be found. On the other hand, the studies of BD characteristics of these dielectrics are strongly complicated by their double-layer structure due to the presence of inevitable lower- $k$  interface region at Si, i.e. the film has to be treated as stacked layer. So the complexity of high- $k$  materials defines entirely different mechanism causing BD from that in  $\text{SiO}_2$ , and the models developed for  $\text{SiO}_2$  can not be easily implemented. Moreover, field acceleration models are not applicable and it is more correct to use voltage acceleration approaches.

This work addresses issues such as hard BD, the effect of traps on the breakdown characteristics of both pure and doped  $\text{Ta}_2\text{O}_5$ , and the role of the dopant in these processes.  $\text{Ta}_2\text{O}_5$  is a leading high- $k$  dielectric for storage capacitors in dynamic memories. The doping of  $\text{Ta}_2\text{O}_5$  results in improved stack characteristics [5], [8], [12]-[14] allowing further scaling of equivalent oxide thickness (EOT), and in terms of leakage current and dielectric constant aspects the effect of various dopants is somehow clarified. Our previous studies [5], [8], [13], [14] revealed that the doping improves the overall electrical characteristics of the stack through a compensation of oxygen vacancies in  $\text{Ta}_2\text{O}_5$ . At the same time, the dopant initiates the generation of specific traps which affect the current behaviour during electrical stress.

## 2. EXPERIMENTAL PROCEDURE

P-type (100) Si wafers were used as substrates and MIS capacitors with pure, with lightly Al-doped  $\text{Ta}_2\text{O}_5$ , and with Hf-doped  $\text{Ta}_2\text{O}_5$  were prepared.  $\text{Ta}_2\text{O}_5$  film was depos-

ited by reactive sputtering of Ta target in an Ar + 10% O<sub>2</sub> mixture. Details on the deposition conditions are given elsewhere [15]. Al-doped Ta<sub>2</sub>O<sub>5</sub> layers were deposited by rf sputtering of Al/Ta alloy target (5 at.% Al content) in Ar + 10 % O<sub>2</sub> atmosphere following the procedure described previously [16]. Hf-doped Ta<sub>2</sub>O<sub>5</sub> was deposited on the Si as it is described in [17], (~ 0.7 nm thick Hf layer was deposited by sputtering on the top of ~7.5 nm rf sputtered Ta<sub>2</sub>O<sub>5</sub>). Post-deposition annealing was carried out at 400°C in pure N<sub>2</sub> ambient in order to mix Hf and Ta<sub>2</sub>O<sub>5</sub> and to ensure diffusion of Hf into te matrix of Ta<sub>2</sub>O<sub>5</sub>.

The films are amorphous according to X-ray diffraction analysis, and the basic physical and electrical properties of the doped films have been reported elsewhere [5], [13], [17], [18]. The results suggest that the doping tends to create shallow traps which mediate the charge transport in doped films thereby changing the dominant conduction mechanism as compared to the pure Ta<sub>2</sub>O<sub>5</sub>. The overall thickness of the films including interfacial layer determined ellipsometrically is 7-8 nm (the interfacial layer is SiO<sub>2</sub>-like one with a thickness of ~ 1 nm). The capacitors with back side electrode of 300 nm evaporated Al and sputtered W top electrode with gate area of 1x10<sup>-4</sup> cm<sup>2</sup> were defined by photolithography. W was chosen as a metal gate on account of its good chemical stability on Ta<sub>2</sub>O<sub>5</sub>-based films. A post-metallization annealing was performed in a forming gas at 450°C for 1 h. The electrical stress response of the capacitors is investigated by monitoring the BD statistics of leakage current dependence on time. Constant voltage stress (CVS) with a negative bias on the top electrode (gate injection) is applied to study the TDDB characteristics. The applied voltage is kept constant and the change in the current density J is monitored as a function of the stress time,  $t_s$ . A breakdown event is defined as a sudden and distinct increase in the current; the  $t_s$  increases until the current exceeds 1 μA. Hard breakdown (HBD) is taken as a BD criterion. Every  $t_{bd}$  – BD data point at 63.2 % value ( $t_{63}$ ) is obtained from Weibull distribution of 40 samples.

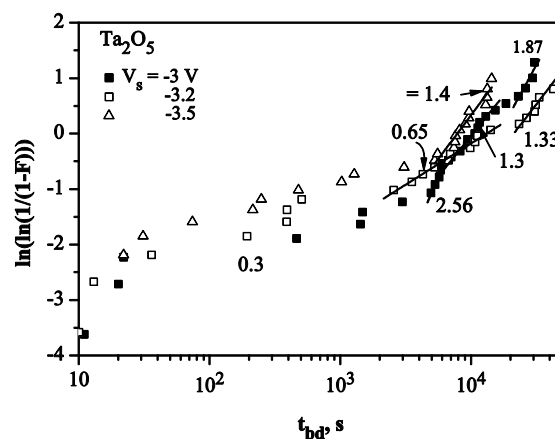
### 3. RESULTS AND DISCUSSIONS

The time-to-breakdown  $t_{bd}$  is statistically distributed and is described by the Weibull distribution:  $F(t) = 1 - \exp(-t/t_{63})^\beta$ , where F is the function of cumulative fraction of failed devices;  $t$  denotes the variable of time-to-breakdown;  $t_{63}$  is the characteristic time for 63.2 % of the devices to fail,  $\beta$  is the shape factor of Weibull distribution (Weibull slope). Plotting the Weibull function in the form  $W = \ln[-\ln(1-F(tbd))]$  vs.  $\ln t_{bd}$  yields a straight line with a slope  $\beta$ , which entity is a measure for the BD behaviour. The dielectric quality reflects on the value of  $\beta$ : for nearly perfect SiO<sub>2</sub>,  $\beta > 1$  in the intrinsic region where the breakdown is an inherent property of SiO<sub>2</sub>. The high- $k$  dielectrics possess large density of electrically active defects which enhance the dielectric failure. In dependence on the nature of defects, their density and stress conditions including the measurement temperature, the slope of Weibull distribution could be lower or higher than 1.

#### 3.1. Pure Ta2O5 films [19]

The capacitors are stressed at voltages between -3 and -3.5 V until BD, with 40 samples for each voltage, allowing for an accurate extraction of the statistical parameters, (for stress voltages more negative than about -3.7 V, the time-to-breakdown is extremely short and accurate characteristics can not be obtained);  $t_{bd}$  is defined as hard breakdown occurs. There

is a correlation between  $t_{bd}$  and leakage current  $J_0$  in the initial samples: devices with low  $J_0$  break later than those with high one. Fig. 1 shows the Weibull distributions of capacitors under various  $V_s$  at room temperature. The distribution is with a slope of  $\beta$  at high percentiles higher than that at low ones. This is due to the coexistence of extrinsic and intrinsic modes: the tendency of  $\beta$  to be smaller at low percentiles is typically observed for the extrinsic/intrinsic BD modes. There is no visible dependence the steeper slope to be observed at higher  $V_s$  implying that the BD is not accelerated by  $V_s$ . The explanation could be sought in the presence of process-induced defects in  $Ta_2O_5$  which hamper the manifestation of well pronounced stress-voltage dependence of  $\beta$ . In the short time range  $\beta$  is  $\sim 0.3$ , in the long time one the slope is steep, i.e. a separate detection of early and long-time BD is visible, and the region with very small Weibull slope proceeds to a higher-slope distribution. The early BDs are considered as a result of weak localized regions formed after critical number of traps. When discussing Weibull distribution of high- $k$  stacks, a number of factors must be considered. The Weibull slopes are limited by the structure and chemical composition of both the interface layer and the bulk high- $k$  layer. Therefore, the values of  $\beta$  should be defined by mechanisms of degradation of these two sublayers. The effect of interface layer is essential due to the high electric field across this layer. This layer is physically thinner than the  $Ta_2O_5$  and is exposed to high electric field due to its lower dielectric constant. A part of the voltage drop across the interface region during the stress will drop across the  $Ta_2O_5$  after interface degradation. Considering the double-layer structure of the high- $k$  film the early breakdowns are assigned to extrinsic BDs, whereas *the breakdowns caused by the interface layer itself could be defined as intrinsic BDs*, since the presence of the interface layer is an inherent property of high- $k$  stack. The defined by this way intrinsic BD explicitly includes the effect of interface layer and its BD. So all early breakdowns with  $t_{bd}$  from 10 to  $\sim 10^3$  s, are related to extrinsic BDs and they characterize weak spots in the films. Their amount is small and the main fraction of failures are actually with  $t_{bd} \sim 10^3$ - $5 \times 10^4$  s. The observed difference in  $\beta$  values for early and intrinsic BDs can not be due only to statistical effects and seems to be additionally related to the strong effect of interface lower- $k$  layer at the Si which not only triggers BD but influences its Weibull distribution.



**Fig. 1** TDDB Weibull distribution of  $Ta_2O_5$  capacitors under different stress voltages. The values of Weibull slope are given

Ultimately, the low Weibull slope values typically observed in high- $k$  stacks even in the regions considered as intrinsic BD, are assigned to low critical defect density needed to cause BDs.  $\beta$  value, however, has been involved to explain single SiO<sub>2</sub> layer BD rather than double-layer high- $k$  stack. The observed quite lower values of  $\beta$  in the region of long  $t_{bd}$  for Ta<sub>2</sub>O<sub>5</sub> than SiO<sub>2</sub> films with the same physical thickness means that the percolation model developed for SiO<sub>2</sub> does not seem to apply to Ta<sub>2</sub>O<sub>5</sub> layers. So we can anticipate that  $\beta$  values even in this  $t_{bd}$  region correspond to the fact that the BD distribution is basically controlled by process-induced defects forming weak spots which are randomly distributed. From the other hand, if the BD is caused by stress-induced traps and the percolation model for SiO<sub>2</sub> is valid for high- $k$  Ta<sub>2</sub>O<sub>5</sub>,  $\beta$  is proportional to the number of traps which form a conducting path. Thus, the lower  $\beta$  values of Ta<sub>2</sub>O<sub>5</sub> mean lower number of stress-induced traps (as compared with SiO<sub>2</sub> with the same physical thickness) which constitute the BD [20]. This is only possible if the trap radius is significantly larger than in SiO<sub>2</sub>. Thus, the use of percolation concept directly results in low critical trap density with very large trap radius at BD. Therefore, a reason for the relatively low Weibull slopes in the range of  $t_{bd} \sim 10^4$  s could be a large effective radius of defects involved in BD events [21]. As is seen, the percolation approach for Ta<sub>2</sub>O<sub>5</sub> stacks can be accepted only if we assume the small critical defect density at breakdown combined with large spacing between defects where tunneling of trapped electrons becomes probable, and this leads to a small Weibull slope. If we consider, however, EOT of Ta<sub>2</sub>O<sub>5</sub> (2.2 nm) instead of its physical thickness,  $\beta$  values in the range of long  $t_{bd}$  ( $> 2 \times 10^3$  s) are generally in accordance with those of  $\sim 2$  nm SiO<sub>2</sub> layer ( $\beta \sim 1$  [22]), and consequently the percolation model could be used for Ta<sub>2</sub>O<sub>5</sub>. To complete the discussion of TDDB characteristics we examine the curves whose  $t_{bd}$  is close to  $10^4$  s, (Fig. 2). If the percolation model is valid the curves should be shifted to the left with increasing  $|V_s|$ , i.e.  $V_s$  acts as an acceleration factor for BD. This is fulfilled for  $V_s = -3.2$  and  $-3.5$  V. The position of the data for  $V_s = -3$  V does not correspond to the percolation concept. This means that the mechanism of degradation in this case is quite specific and the extrinsic explanation seems reasonable. We speculate that the TDDB characteristics with  $t_{bd} \sim 10^4$  s, at  $V_s = -3$  V are attributed to the activation of pre-existing traps during the stress and these traps presumably dictate the current behaviour. The physical nature of the involved traps are oxygen vacancies (typically ob-

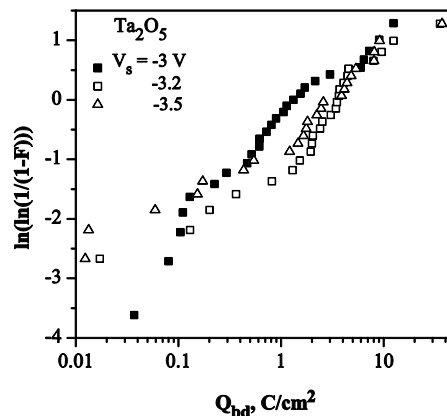


Fig. 2 Charge-to-breakdown distribution for Ta<sub>2</sub>O<sub>5</sub> devices under CVS

served in Ta<sub>2</sub>O<sub>5</sub>) and gate-deposition induced defects during sputtering of W electrode. In opposite, new traps generation dominates at higher  $V_s$ , the dielectric quickly goes on to HBD, and  $V_s$  is an acceleration factor for BD. The generation of new defects (a formation of the percolation path) could additionally depend on the local field produced by the initial traps. Accordingly, this would affect the way of formation of the conductive path(s). So to all appearance both the pre-existing traps and the new traps generated affect the region with long  $t_{bd}$ : at  $V_s = -3$  V dominate extrinsic failure modes and at -3.2; -3.5 V percolation concept appears to more accurately describe the results.

Charge to breakdown  $Q_{bd}$  is a more effective tool than  $t_{bd}$  for evaluating the generated defect density per injected electron. The Weibull distribution of  $Q_{bd}$  is shown in Fig. 3. The  $Q_{bd}$  is obtained by the integral of the leakage current over time ( $Q_{bd} = \int_0^{t_{bd}} J dt$ ), so all current bursts are considered. The  $Q_{bd}$  values are 1.3 C/cm<sup>2</sup> at 63.2 % failure of Ta<sub>2</sub>O<sub>5</sub> for  $V_s = -3$  V and comparable ones (3.6 C/cm<sup>2</sup>) for the stress biases of -3.2 and -3.5 V, (63.2% are the percents corresponding to  $\ln[-\ln(1-F)] = 0$ ). Accordingly, a higher  $Q_{bd}$  is obtained at larger stress voltages. The higher  $Q_{bd}$  means lower defect generation rate. The behaviour of  $Q_{bd}$  confirms that the mechanism of degradation for  $V_s = -3$  V is different from that for -3.2 and -3.5 V. In the former case the trapping at pre-existing traps dominates, entails the appropriate conditions for BD, and the effect of  $V_s$  as a factor for generation of new traps can not be observed. There is no relative change of the current at BD which could constitute well defined  $Q_{bd}$ . In terms of  $t_{bd}$  this means that the time needed to create the local increase of current resulting in BD event has not a certain value. The reason is the dominant effect of weak spots. This makes the use of the current variation just before BD as a measure for evaluating the rate  $R_{gen}$  at which defects are generated and the defect density  $N_{bd}$  required to cause BD, strongly unreliable. The long BD times and the evidence of weak spots could be interpreted with an impact of both the process-induced defects and the interfacial layer on the BD behaviour. Therefore, the results demonstrate that *the picture for BD in Ta<sub>2</sub>O<sub>5</sub>-based capacitors is quite complicated. The origin of BD is not the same as that in single SiO<sub>2</sub> which is expectable, having in mind the double-layer structure of Ta<sub>2</sub>O<sub>5</sub> films on Si. The role of weak spots is essential in BD behaviour, but*

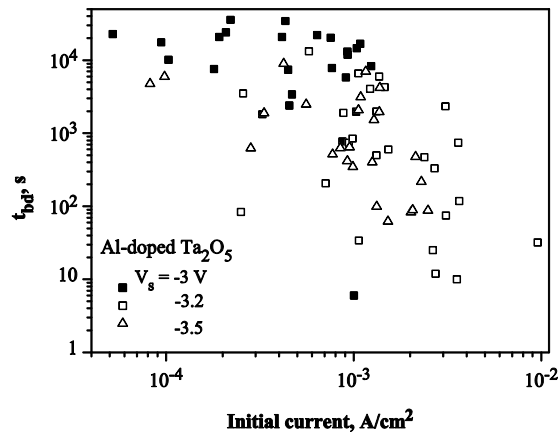
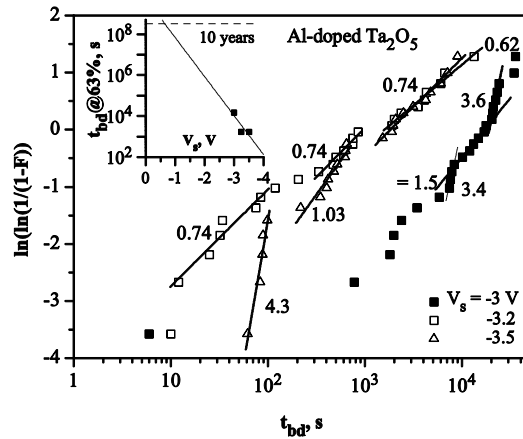


Fig. 3  $t_{bd}$  vs. initial current in Al-doped Ta<sub>2</sub>O<sub>5</sub> devices for different CVS levels

the other factors such as the quality of interface region and its thickness homogeneity could not be discarded as a contributing factor especially in the so called "intrinsic" BD.

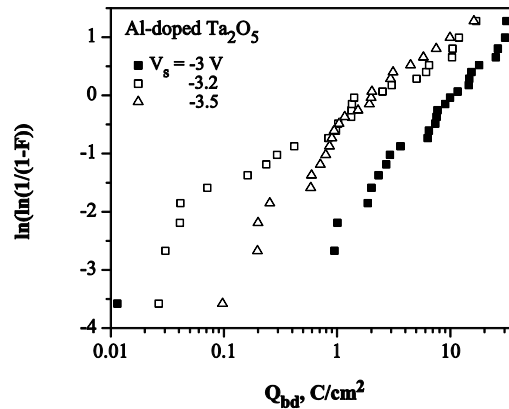


**Fig. 4** Weibull distributions of  $t_{bd}$  at various stressing voltages for Al-doped  $Ta_2O_5$  capacitors. The inset shows the  $t_{63}$  values vs. stress voltage

### 3.2. Al-doped Ta<sub>2</sub>O<sub>5</sub> [19]

The BD occurs a little faster for the doped layers ( $EOT = 1.5$  nm) as compared to the  $Ta_2O_5$ , i.e. it emerges that the presence of Al speeds up the breakdown.  $t_{bd}$  depends on both the bias stress and the initial current of capacitors (Fig. 3):  $t_{bd}$  is smaller at larger  $V_s$  and larger initial current. The observed multimodal BD distribution (Fig. 4) with varying slope for given  $V_s$  suggests the presence of extrinsic failure caused by processing defects. The slope of the curves for  $t_{bd}$  over  $\sim 10^3$  s is due to intrinsic BD in terms of high- $k$  stacks defined above. Quite high  $\beta$  values of doped films for  $t_{bd} \sim 10$ - $10^3$  s, (up to  $\beta = 4.3$ , against  $\beta \sim 0.3$  for undoped films) indicate that Al effectively reduces some of the defects in  $Ta_2O_5$ . Nonuniform distribution of weak spots causes a variation in  $t_{bd}$  and leads to variation in  $Q_{bd}$  as well (Fig. 5).  $Q_{bd}$  at 63.2 % failure is  $10.5$  C/cm<sup>2</sup> when  $V_s = -3$  V, and  $Q_{bd}$  values are equal ( $2.1$  C/cm<sup>2</sup>) after stressing with  $-3.2$  and  $-3.5$  V. Compared to the  $Ta_2O_5$ , the  $Q_{bd}$  values are quite larger in the former and smaller in the latter case, respectively.

The curve for  $V_s = -3$  V for doped films shifts to the right resulting in a reasonable behaviour of  $Q_{bd}$  as a function of  $V_s$ , i.e.  $Q_{bd}$  is reduced at higher stress voltage which means higher defect generation rate at higher  $|V_s|$  for the doped sample if the percolation model is valid. This behaviour is in contrast to the case of  $Ta_2O_5$  and implies that Al significantly affects the defects defining the BD at  $V_s$  of  $-3$  V. Compared with the  $Ta_2O_5$  the  $Q_{bd}$  and TDDB characteristics for the doped samples are improved which means lower charge trapping. The maximum voltage to guarantee a 10 years lifetime is found to be  $-0.65$  V for Al-doped  $Ta_2O_5$  (inset of Fig. 4,  $t_{bd}$  values at 63.2% are plotted vs.  $V_s$ ). It was not possible to extract accurately the operating voltage for 10 years lifetime for  $Ta_2O_5$  due to the very small variation of  $t_{63}$  with  $V_s$  (Fig. 1). The doped capacitors are also better than  $Ta_2O_5$  ones in life-time scaling due to the larger  $\beta$  value of the former.

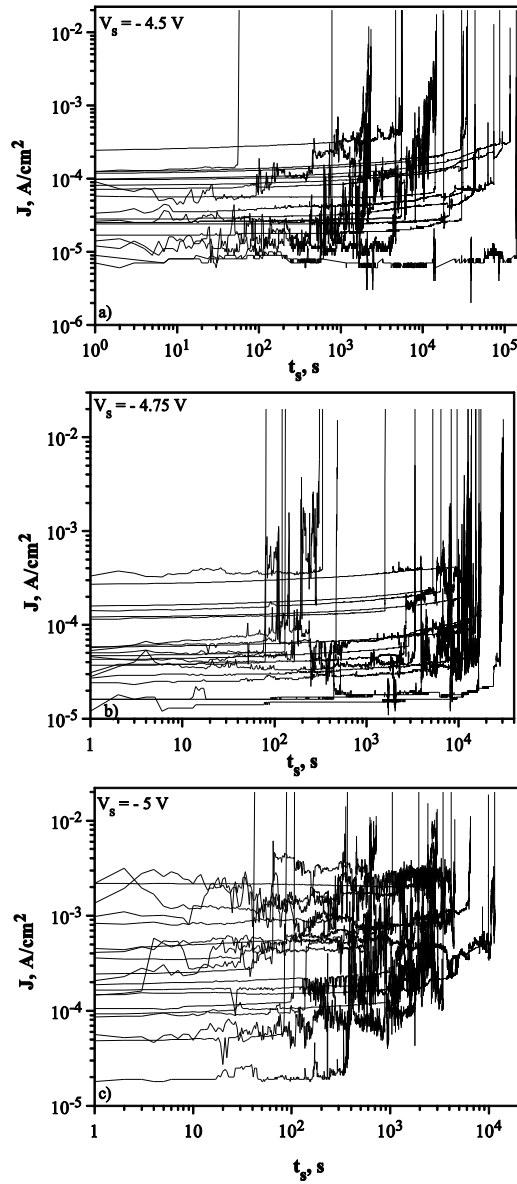


**Fig. 5** Weibull distributions of  $Q_{bd}$  for Al-doped  $Ta_2O_5$  under various  $V_s$

All those results indicate that the BD in both *undoped and Al-doped  $Ta_2O_5$*  is affected by the process-induced defects and it cannot completely be explained only by the percolation model. The presence of defects forms weak spots which dictate the BD. This is compounded by the double-layer structure of the high- $k$  stack and the impact of the interface region. As a result, *the value of Weibull slope is a total product of these two effects and can not exhibit by default any clear trend*. When the effect of interfacial layer dominates over the effect of defects the BD is attributed to intrinsic BD of the stack. This also means that the dominant mechanism of degradation is embodied in the density of pre-existing defects in high- $k$  dielectric. In this context two different mechanisms of BD emerge for pure and Al-doped  $Ta_2O_5$ : i) the pre-existing traps causing weak spots in undoped layers play a key role in electrical degradation; they mainly constitute the BD and extrinsic mechanism of degradation is prevalent, especially at low voltage; at high  $V_s$ , stress-generated new traps appear to initiate the BD. ii) Al doping reduces defects in  $Ta_2O_5$  resulting in the domination of intrinsic failures in the breakdown statistics, and the mechanism of degradation is more adequate (but not completely!) to the percolation concept. The type of defects, however, are different in two type samples regarding their structural nature, amount and energy location, which factors reflect on the TDDB data and could explain the observed difference between undoped and doped  $Ta_2O_5$ . According to the TEM data, the weak spots are not local physical thinning of the dielectric, (pure and doped one). The SILC-measurements of W-gated  $Ta_2O_5$  capacitors under CVS at gate injection have revealed that the processes of electron trapping at pre-existing traps and positive charge build-up are responsible for the SILC [23]; there was no evidence of stress-induced generation of new traps. Unlike the case of  $Ta_2O_5$ , the SILC in Al-doped  $Ta_2O_5$  is a result of new traps generation [24]. The stress-generated traps act as stepping (transport) sites for injected electrons giving rise to SILC, (i.e. the electrons are trapped and released immediately thereafter). Therefore, considering the previous SILC measurements, it is reasonable to interpret the results in terms that *BD mechanism in pure and lightly Al-doped  $Ta_2O_5$  is intimately related to both pre-existing traps and stress-induced new traps*. In the former case the impact of trapping at pre-existing traps is a little dominant and in the latter one new traps generation mainly defines the BD process. In this sense the BD degradation of doped films is dominated by the Al-induced modification of the  $Ta_2O_5$  network. Regarding microstructural aspects of the traps, wide category of bond



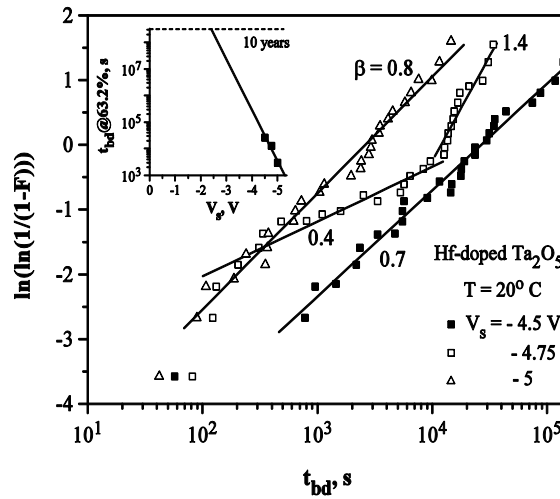
defects (weakened and non-perfect Ta-, Si- and Al-related bonds) and oxygen vacancies represent precursors for BD defects. The bond-breaking is suggested as a general mechanism to describe the electrical breakdown behaviour.



**Fig. 6** TDDB characteristics of Hf-doped Ta<sub>2</sub>O<sub>5</sub> stacks stressed at room temperature with  $V_s$  of - 4.5 (a); - 4.75 V (b); - 5 V (c)

### 3.3 Hf-doped Ta<sub>2</sub>O<sub>5</sub> [25]

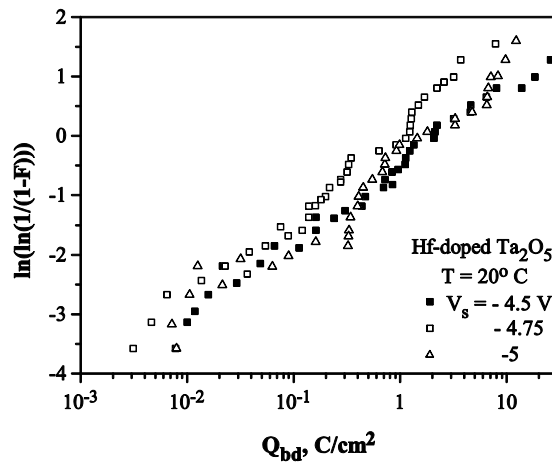
Fig. 6 shows the leakage current density-stress time ( $J$ - $t_s$ ) characteristics under various  $V_s$  at room temperature. The curves demonstrate relatively constant current up to BD, similarly to the case of pure and lightly Al-doped Ta<sub>2</sub>O<sub>5</sub>. With similar values of  $J_0$ , pure Ta<sub>2</sub>O<sub>5</sub> stacks reveal shorter time-to-breakdown. The longer  $t_{bd}$  in the doped films may be attributed to the compensation of oxygen vacancies by Hf in Ta<sub>2</sub>O<sub>5</sub> accompanied by the observed improvement of the stack parameters. The structural nature of the weak spots is regions with a concentration of bond defects (weakened and non-perfect Ta-; Si-; and Hf-related bonds).



**Fig. 7** Weibull plots of Hf-doped Ta<sub>2</sub>O<sub>5</sub> under different  $V_s$ . The inset shows  $t_{63}$  values vs. stress voltage

Weibull plots for various  $V_s$  are given in Fig.7. The distribution is with similar slopes for the lowest and the largest  $V_s$ . In the short time range ( $10^2$ - $10^4$  s) low  $\beta$  value (0.4) corresponds to the early BDs; the slope is steep (1.4) in the longer time range. Generally, the Weibull distribution of TDDB differs from that of Ta<sub>2</sub>O<sub>5</sub> where very clear separation of early and long-time BDs was detected with different  $\beta$  values. The lack of well pronounced region of early breakdowns for the doped samples implies that they have less process-induced defects. Considering the bi-layer structure of the films and inevitable difficulties in correct estimation of the applied fields in the interface layer and in the high- $k$  layer we again discuss the BD characteristics here in terms of ones controlled by applied voltage rather than the applied field. We will focus on the long-time BD (more than  $\sim 10^4$  s) as the most relevant to the intrinsic electrical degradation of high- $k$  films and where the influence of the interfacial region on the breakdown phenomenon is explicitly included. The detected lower  $\beta$  values (0.7-1.4) than the value for SiO<sub>2</sub> with the same physical thickness indicates that the percolation model does not seem to apply to Hf-doped Ta<sub>2</sub>O<sub>5</sub> films, and the extrinsic explanation of HBD seems more likely. The percolation model can be accepted only if we assume low number of defects involved in BD with large spacing between defects where tunneling of trapped electrons becomes probable and causes low value of Weibull slope. (Note that the charge trapping could be everywhere in

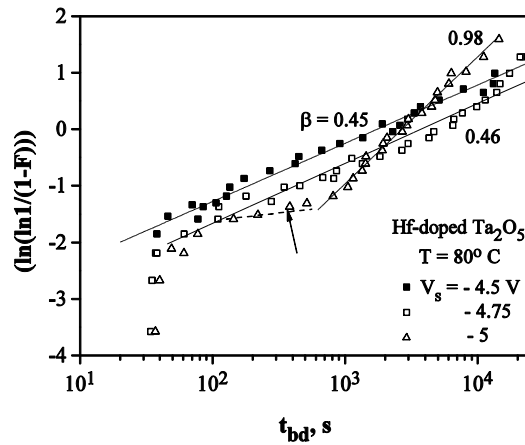
the stack, i.e. in the bulk high- $k$  as well as within the interface layer; the trapping in the interface layer could not only trigger breakdown of the stack but influences its Weibull distribution.) If we consider EOT of Hf-doped Ta<sub>2</sub>O<sub>5</sub>, however, instead of its physical thickness,  $\beta$  values in the range of  $t_{bd} > 10^4$  s are in accordance with those of  $\sim 2$  nm SiO<sub>2</sub> layer, and consequently the percolation model for the stressed Hf-doped stack can be used. An additional support for the percolation approach is the fact that the curves are shifted to the left with increasing  $|V_s|$ , i.e. the stressing voltage acts indeed as an acceleration factor for BD, and the generation of new defects affects the HBD (Fig. 7). This behavior of the Weibull distribution also indicates improved BD characteristics of the Hf-doped Ta<sub>2</sub>O<sub>5</sub> with respect to undoped one. The maximum projected voltage for a 10 years lifetime is found to be -2.4 V, or equivalently 3 MV/cm, (inset of Fig. 7). This parameter is significantly larger than that of SiO<sub>2</sub> [26] and Al-doped Ta<sub>2</sub>O<sub>5</sub> with the same EOT. Therefore, lifetime scaling of Hf-doped Ta<sub>2</sub>O<sub>5</sub> is definitely better than both pure and Al-doped Ta<sub>2</sub>O<sub>5</sub>, and might be better than SiO<sub>2</sub> when compared the same EOT value.



**Fig. 8**  $Q_{bd}$  distribution for Hf-doped Ta<sub>2</sub>O<sub>5</sub> under  $V_s$  from - 4.5 to - 5 V

The Weibull distribution of charge-to-breakdown is shown in Fig. 8. The curves corresponding to  $V_s = - 4.5$  and  $- 4.75$  V are as if in accordance with the percolation model (the curve with more negative  $V_s$  is shifted to the left). The curves for the lowest and highest  $|V_s|$ , however, are indistinguishable. Since the  $Q_{bd}$  variation does not appear to have a definite behavior as a function of  $V_s$ , it is not correct to consider  $V_s$  as an acceleration factor for breakdown. The trapping at pre-existing traps at  $V_s = - 5$  V obviously dominates and the effect of  $V_s$  as a factor for generation of new traps can not completely be revealed. The observed phenomenon corresponds to the long-term  $t_{bd}$  where the role of interfacial region in the BD events is essential. Thus the effect of charge trapping in the interface layer itself could be dominant on the BD process. We found linear relationship between  $Q_{bd}$  and  $t_{bd}$  implying that the average leakage current at BD has narrow distribution and the distribution of  $t_{bd}$  is caused by large size defects. Similarly to the case of pure and Al-doped Ta<sub>2</sub>O<sub>5</sub>, SILC at breakdown for Hf-doped films can not be used as a measure of increasing defect density during stressing due to the strong effect of weak spots. As it

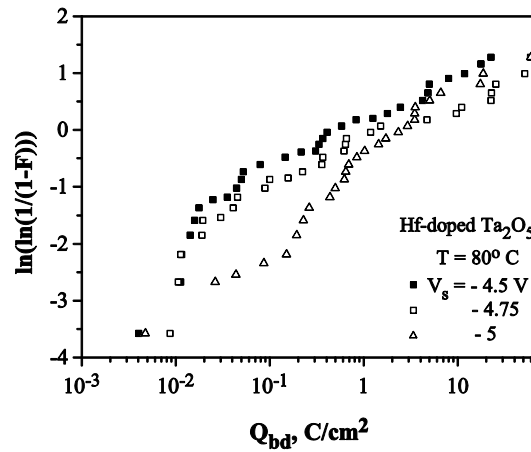
was reported earlier [6] electron trapping at pre-existing sites in high- $k$  dielectrics can be the first stage of stress-induced trap generation process. The strong trapping phenomena, however, results in very complex relation between the BD and SILC, completely different from that for SiO<sub>2</sub>. The evidence for the presence of weak spots in these experiments in combination with long  $t_{bd}$  indicates that both the process-induced defects and the interfacial layer affect HBD. The action of these factors leads to specific Weibull distribution and to unpredictable dependence of  $t_{bd}$  and  $Q_{bd}$  on  $V_s$ . The introduction of Hf into the matrix of Ta<sub>2</sub>O<sub>5</sub> generates shallow traps which are responsible for the domination of the tunneling processes in the conductivity in Hf-doped Ta<sub>2</sub>O<sub>5</sub> [13], [17]. These traps in contrast to the expectation do not speed up breakdown, and  $t_{bd}$  for the doped film is even larger than in the pure Ta<sub>2</sub>O<sub>5</sub> under the same stressing voltages. Presumably, the presence of Hf into Ta<sub>2</sub>O<sub>5</sub> ultimately improves the BD characteristics and modifies the dielectric BD mechanism(s) manifesting itself as differences in the breakdown statistics of undoped and doped films.



**Fig. 9** TDDDB Weibull distribution of Hf-doped Ta<sub>2</sub>O<sub>5</sub> devices under various stress voltages at 80°C

The high temperature stressing at 80°C modifies Weibull distribution to some extent - the breakdown occurs faster than that at room temperature, but there is no visible temperature dependence of the shape factor  $\beta$  (Fig.9). The combined effect of high T and stressing voltage on the Weibull slopes does not exhibit unambiguously trend; the percolation model for BD at 80°C is not applicable for all  $V_s$ . The effect of  $V_s$  at high temperature on the BD has no simple explanation, but nevertheless there are indications for TDDDB reliability degradation: the stress at high temperature worsens the characteristics, stimulates charge trapping with evidence that large size defects dictate the breakdown; TDDDB lifetime degrades, and it is not possible to extract the operating voltage for 10 years lifetime. The results indicate for some redistribution of defects at high temperature for both doped and undoped samples. For example, the presence of regions with different slopes at a certain  $V_s$  and a single plateau between them (Fig. 9, the region indicated by arrow) implies for the activation of different type traps and their dominance in the conductivity of Hf-doped films. The nature of defects is not the same for the pure and doped Ta<sub>2</sub>O<sub>5</sub> which

manifests as differences in the corresponding Weibull distributions. The stressing at high temperatures weakens some bonds and results in enhanced trap generation. One can not also exclude the effect of hydrogen-related particles on the degradation at elevated temperatures. Hydrogen is ubiquitous in the high- $k$  system and can be everywhere in the stack. During stressing at high temperature liberated protons at the interface with Si could be a key factor for enhanced charge trapping at strained bonds in the interface layer. Being liberated hydrogen readily diffuses in the high- $k$  layer resulting in generation of positively charged centers or neutral trapping centers.



**Fig. 10** Weibull distributions of  $Q_{bd}$  at 80°C for Hf-doped Ta<sub>2</sub>O<sub>5</sub> under different  $V_s$

The charge to breakdown for the doped films at 80°C does not reduce at large stress voltage and the  $Q_{bd}$  data can not be interpreted within the percolation model: instead, the  $Q_{bd}$  curves are shifted to the right with increasing  $|V_s|$  (Fig. 10). The observed higher  $Q_{bd}$  at largest  $V_s$  generally means a lower defect generation rate. The shallow traps which are an attribute of Hf-doped Ta<sub>2</sub>O<sub>5</sub> could be responsible for this phenomenon.

Referring to the conductivity and SILC measurements [5], [17], we note that distinctly different conduction mechanisms control the leakage current in pure and Hf-doped Ta<sub>2</sub>O<sub>5</sub> which are defined by the traps participating in the current transport. The oxygen vacancies are the main defects in Ta<sub>2</sub>O<sub>5</sub> which constitute the domination of Poole-Frenkel mechanism. The conduction in Hf-doped Ta<sub>2</sub>O<sub>5</sub> is performed through shallow traps resulting in domination of tunneling processes in the conductivity. The consequence of the transport through shallow traps is considerably weaker T-dependence of the current. One possible scenario of the HBD mechanism in Hf-doped films includes breaking at strong stress of weakened from the charge trapping bonds; this results in new traps generation (i.e. percolation model is valid) which quickly leads to HBD. The broken bonds could be precursor breakdown defects. Unlike the case of Ta<sub>2</sub>O<sub>5</sub> where the BD mechanism is dominated by the trapping at the pre-existing traps (deep traps attributed to oxygen vacancies; the vacancy can trap one or two electrons), new traps generation at strong stress becomes more effective for Hf-doped Ta<sub>2</sub>O<sub>5</sub>, and affects its final BD. The HBD in the doped Ta<sub>2</sub>O<sub>5</sub> is presumably dictated by the Hf-induced shallow traps instead of deep oxygen vacancies.

This indicates that the Hf-induced modification of Ta<sub>2</sub>O<sub>5</sub> network governs the TDDB behavior manifesting as stress-generation of new traps. These in-fabricated traps cause weak spots in the doped films and account for the early breakdowns. At high levels of stress, which are used to examine TDDB characteristics, the weakened by the trapping bonds break and play the role of a nucleation center for the generation of defect clusters that can form a conduction link through the dielectric until HBD occurs, i.e. stress-generated new traps process becomes significant and the mechanism of BD is more adequate to the percolation approach. In this context electron trapping at pre-existing traps is a concomitant process for BD. Therefore, *the BD mechanism in Hf-doped Ta<sub>2</sub>O<sub>5</sub> stack is associated with both pre-existing shallow traps and stress-induced new traps*. In the final stage of BD the last process is dominant. These processes are compounded with the impact of the interface region and its breakdown. When the degradation of interface layer dominates over the effect of the initial defects in the high-*k* layer the BD is designed as an inherent property of high-*k* stack, i.e. as its intrinsic BD. In the term of  $t_{bd}$  this corresponds to long  $t_{bd}$  ( $> 10^4$  s) where the high applied field on the interface region quickly provokes breakdown. So the combined effect of the initial defects, the stress-induced new traps and the effect of the interface region itself constitutes the Weibull distribution of dielectric BD in high-*k* systems.

#### 4. CONCLUSION

The picture of the BD in Ta<sub>2</sub>O<sub>5</sub>-based films on Si is quite complicated and different than that of SiO<sub>2</sub> with the same physical thickness. The pre-existing traps which provoke weak spots with non-uniform distribution are responsible for the early BDs. Their effect is strong enough, even in the range of long time-to-breakdown where the impact of inevitable lower-*k* interface layer at Si can not be neglected. Since the presence of this layer is inherent property of high-*k* film on Si, it is reasonable to designate the failure of the interface layer as an intrinsic mode of the BD of high-*k* stack. Thus, the behaviour of Weibull distribution of time-to-breakdown for pure and doped Ta<sub>2</sub>O<sub>5</sub> could be interpreted as co-existence of extrinsic and intrinsic BD modes. Both the pre-existing traps and the new traps generated affect the region with long time-to-breakdown. The domination of one of these competitive effects dictates the mechanism of degradation: the effect of process-induced defects dominates in Ta<sub>2</sub>O<sub>5</sub> and the BD can not completely be explained by the percolation approach: the trapping at pre-existing traps facilitates the breakdown and the effect of stressing voltage as a factor for generation of new traps can not be well observed. The doping reduces defects in Ta<sub>2</sub>O<sub>5</sub> which manifests in larger Weibull slopes and better TDDB characteristics of the doped samples: the generation of new traps prevails over the charge trapping and the mechanism of breakdown is more adequate to the percolation model. Accordingly, the BD degradation in doped Ta<sub>2</sub>O<sub>5</sub> is dominated by the dopant-induced modification of Ta<sub>2</sub>O<sub>5</sub> network indicating that the dopant element in high-*k* dielectrics could be a primary factor in reliability issues.

Lifetime scaling of the doped films is better than Ta<sub>2</sub>O<sub>5</sub> and might be better than SiO<sub>2</sub> with the same equivalent oxide thickness. Trapping phenomena, however, are still observed in the doped films which complicate TDDB characteristics; in some cases the effect of stressing voltage as an acceleration factor for electrical degradation is obscure.

The influence of pre-existing traps is compounded with the effect of high applied field in the interface layer and of stress-induced new traps generation in the range of long-time-to-breakdown. The action of these factors leads to specific Weibull distribution of the high- $k$  samples and to unpredictable dependence of time and charge to breakdown on the stressing voltage. *The peculiarities of the TDDB characteristics of the pure and doped Ta<sub>2</sub>O<sub>5</sub> are assumed to be due to the inherently different process-induced defects* (deep levels attributed to oxygen vacancies in the former and shallow traps in the latter case). The trapping at pre-existing traps is the starting point for the electrical failure in both stacks. Further, the trapping at deep traps remains dominant for Ta<sub>2</sub>O<sub>5</sub> up to final stage of breakdown. The HBD in the doped Ta<sub>2</sub>O<sub>5</sub> is defined by the dopant-induced shallow traps: at strong stress the weakened by trapping bonds break and this vastly facilitates new traps generation, and the mechanism of degradation is more adequate to the percolation model. The stressing at high temperatures additionally weakens some bonds and enhances trap generation. When the interface region experience extremely high applied filed, the degradation of interface layer could prevail over the defect-related processes in the high- $k$  layer. Therefore, *three competitive factors constitute the Weibull distribution of dielectric BD in high- $k$  systems*: initial defects, stress-induced new traps and the degradation of the interface layer itself. The domination of one of them depends on the stress conditions, nature of pre-existing traps and the interface layer parameters. A subtle problem for high- $k$  stacks is that the Weibull slope is a total product of these three entities and consequently it can not exhibit by default any clear trend and well pronounced relationship with stressing voltage and stressing temperature. This is indeed observed if one tries to find regularity in  $\beta$  variation with  $V_s$  and T. This implicitly suggests that the simple comparison only of Weibull slopes of high- $k$  stack with the corresponding values of SiO<sub>2</sub> is not the proper approach to interpret the TDDB characteristics of high- $k$  films and could result in inaccuracies in reliability prediction.

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