MATHEMATICAL MODELING OF MATERIALLY NONLINEAR PROBLEMS IN STRUCTURAL ANALYSES (PART I – THEORETICAL FUNDAMENTALS)

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Abstract. Material models describe the way they behave when loaded. The paper presents the development of the model beginning with the simplest linear-elastic and rigid-plastic ones. The basic data in the plasticity theory have been defined, such as criterion and yield (failure) surface, hardening law, plastic yield law and normality condition. Yield criteria of Tresca, Von Mises, Mohr-Coulomb and Drucker-Prager were given separately.

Key words: material models, plasticity, yield criterion, yield surface.

1. Introduction

In order to describe behavior of a material in a suitable way, it is necessary to establish constitutive models (constitutive relations or equations) representing mathematical descriptions of their behavior under external load. The constitutive models are formed by establishing relationship between the stress tensor and strain tensor and represent an idealized, that is, more or less rough description of real behavior. Two most frequently applied basic models, describing material properties are ideally-elastic and ideally-plastic models. Real materials almost never meet the conditions defined by the terms for the given models, but they were, primarily due to their simplicity, indispensable for professional practice. In the recent period, with the advent and development of the plasticity theory, new elasto-plastic models have emerged, much more realistically describing the non-linear characteristics of various granular (friction) materials, such as concrete, soil and rock in civil engineering.

Chronologically speaking, elastic models were the first established models of material behavior. The materials described by this model deform under the action of external forces, and after the termination of the action they assume their original shape. Deformations are reversible, relation of stress and strain is unequivocal, that is, stress state is only a function of the strain state and is independent of the stress path.

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Fig. 1. provides basic types of diagrams of stress-strain in axial load. The simplest relation of stress and strain is linear elastic connection (Fig.1.a) proposed by the Hook's law as early as in 17th century, which is at the same time the most frequently applied one. Deformation is reversible, that is, the loading and unloading diagrams are identical. For the majority of materials in practice, application of such relation can be accepted if the stress state changes are small. Except for the linear-elastic relations, non-linear-elastic relations of stress and strain are frequent (Fig. 1.b)

Further development of constitutive models progressed towards introduction of plastic properties of material into expressions for constitutive equations. In this way, the first elements of the plasticity theory are met in Coulomb who as early as in 1773 proposed the plasticity criterion for soil. The theory of plasticity developed slowly, until the half of 20th century when it started to rapidly develop and engender numerous elasto-plastic models corroborated by experiments. The Elasto-plastic models were first developed for metals, but soon were successfully applied for other materials.

The basic characteristic of elasto-plastic models is introduction of plastic (irreversible) deformations under load. Namely, if a part of deformations after unloading is irreversible, the material, apart from the elastic, also exhibits plastic characteristics which are manifested as plastic deformations. An example of such behavior is displayed in the Fig. 1.c.

In the Fig. 1.d three types of linear elasto-plastic behavior were presented where differences after exceeding the yield point can be observed. Curve 1 represents a case of perfect (ideal) plasticity where after reaching the yield point, deformation increases without an increase of the stress, that is, the material exhibits only plastic deformations, whereas the curves 2 and 3 present the cases of strengthening of weakening of the material. On this basis all elasto-plastic models can be divided in two basic groups:

- ideal elasto-plastic materials (without strengthening or weakening)
- elasto-plastic materials with strengthening or weakening

Fig. 1.e presents a rigid-plastic material which shows only plastic deformations under load.

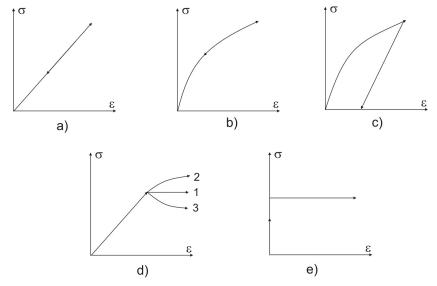


Fig. 1 Idealized behavior of material at axial stress

2. BASIC CONCEPTS OF PLASTICITY THEORY

In order to define the basic concepts encountered in the plasticity theory, behavior of arbitrary elasto-plastic material with strengthening (corresponding to metals), in the conditions of flat deformation, as given in [3]. Namely, two series of sample testing were conducted, the stress σ_c in the first series being kept constant, whereas the stress σ_a , was increased; in the second series the stress σ_c was increased and the stress σ_a was kept constant. Idealized dependency of stress and deformation with the constant value of one stress is given in the Fig.s 2.a and 2.b. It can be observed that with the change of stress σ_y value, apart from the elastic deformations the sample also exhibits the plastic ones. This value of stress is termed the yield stress.

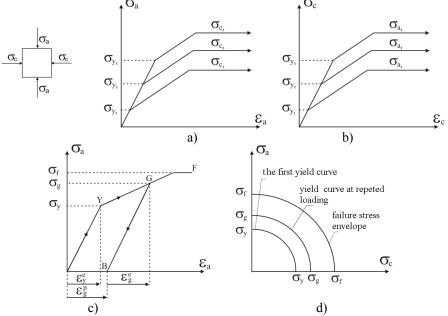


Fig. 2 Behavior of elasto-plastic material with strengthening [3]

If the sample is loaded by stress which is between the yield stress σ_y and failure stress σ_f , that is, by an arbitrary stress σ_g (Fig. 2.c), when unloaded, the sample will retain a part of total deformations in the point $G - \varepsilon_g$. This retained, irreversible part of deformations represents plastic deformations ε_g^p , while reversible, or elastic deformations in the point $G - \varepsilon_g^p$. Under the repeated loading, the sample exhibits plasticity up to the level of stress σ_g when it yields again (but the new yield stress σ_g is higher than the original yield stress σ_y) and it will again exhibit plastic deformations apart from the elastic ones if exposed to further loading. Dependence between the change of yield stress $\sigma_f - \sigma_y$ and plastic deformation from point Y to point $F - \varepsilon_g^p$ is called the law of strengthening or hardening. It defines the way in which the plasticity condition changes in the course of plastic deformation of material.

Elasto-plastic behavior of material can be presented in the diagram $\sigma_a - \sigma_c$ in Fig. 2.d where all the combinations of stresses σ_a and σ_c , where the first yield occurs, are presented by the yield line. All combinations of stresses σ_a and σ_c when material fails are presented by the failure stress envelope. Unloading, then repeated loading in the point G is presented by the line $\sigma_g - \sigma_g$. There is an infinite number of such lines between the first yield line $\sigma_y - \sigma_y$ and failure line $\sigma_f - \sigma_f$ and they form a yield surface between the first yield curve and failure envelope. Therefore if the combination of stresses is such that the point is below the yield line, the specimen will behave elastically, when it is in the yield surface the specimen will exhibit both elastic and plastic deformations, and when it reaches the failure stress envelope, it either experience brittle failure or is plastically deformed without elastic deformations. Outside the stress failure envelope, there are impossible stress states. One should take note the failure stress envelope can be, but need not be, geometrically similar to the failure curves.

Geometrical interpretation of mentioned concepts when applied to construction materials can be presented in the three-dimensional space of main stresses. The failure stress envelope which separates the space of possible and impossible stress states is presented in Fig. 3.a and is also called the failure surface. By its shape, it reminds of a cone whose axis coincides with the main diagonal (where $\sigma_1 = \sigma_2 = \sigma_3$) which is also called the hydrostatic axis. For the materials without cohesion, which cannot receive the tensile stresses (such as the loose soil), the "apex" of the envelope is at the origin, while for the materials with cohesion, it is moved along the hydrostatical axis in its negative direction. The generatrices of the envelope are not entirely straight lines, and the intersection with the octahedron plan, which is perpendicular to the hydrostatic axis, is presented in the Fig. 3.b. The mentioned intersection has three axes of symmetry.

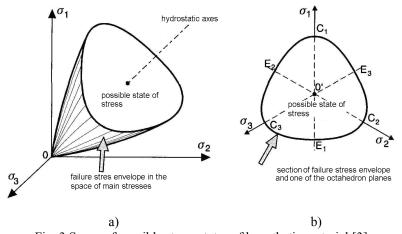


Fig. 3 Space of possible stress states of hypothetic material [3]

In ideally elasto-plastic materials (Fig. 1.d-1) there is no strengthening (hardening) of the material, thus it can be considered that all the yield lines from the Fig. 2 are blended in one and represent the failure envelope. In Fig. 3 the yield surface and failure surface would be practically one surface inside which material behaves in a linear-elastic manner, while the plastic deformations occur when a stress path touches on the failure surface.

For elasto-plastic materials with hardening (Fig. 1.d-2) in the space of main stresses there is an initial yield surface, and then it changes shape and size in the course of plastic deformation. Depending on the way of change of the yield surfaces the material can be with:

- Isotropic (working) hardening
- Kinematic hardening and
- Mixed (anisotropic hardening)

In the materials with isotropic hardening the initial yield surface expands in the main stresses space and remains geometrically similar to the initial one, that is, it does not change shape during the yield of material. In materials with kinematic hardening, the initial yield surface during the plastic deformation changes position in main stresses space, retaining the original size and shape. Such materials exhibit properties expressed by the Bauschinger's effect. By combining of these two ways of hardening, mixed hardening materials are imitated.

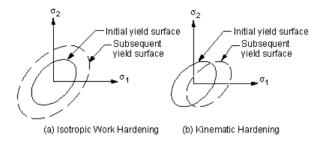


Fig. 4 Types of material hardening [2]

The plastic deformations occur only if the yield surface in one of the mentioned ways is changed or the stress state is in the failure envelope.

The general shape of yield surface can be expressed in the following way:

$$F = F(\{\sigma\}, k) = 0 \tag{1}$$

Where yield surface F is given in dependence of the stress state $\{\sigma\}$ and the hardening parameter k. If F < 0 the stress state is within the yield surface and the material behaves as linear or non-linear elastic one. When F = 0 the stress state is neutral, that is, it is situated on the yield surface for elasto-plastic materials with hardening, while for the ideally plastic materials this is simultaneously the yield surface. For F > 0 a change of the yield surface occurs and begins onset of plastic deformations of elasto-plastic materials with hardening, while this is an impossible stress state for ideally plastic materials.

Regarding that in the plasticity theory no stress-strain linear relations are valid, they cannot be presented in integral but in incremental form, because the relation equations are valid only for one increment. In this way the total increment of deformation can be divided to elastic and plastic components:

$$\{d\varepsilon\} = \{d\varepsilon^e\} + \{d\varepsilon^p\} \tag{2}$$

In isotropic material, the increment of elastic deformation is coaxial with the increment of stress, and the direction and size are determined by the equation:

$$\{d\varepsilon^e\} = [D^e]^{-1}\{d\sigma\} \tag{3}$$

where $[D^e]$ is the matrix of rigidity of material at elastic behavior, and $\{d\sigma\}$ is the increment of stress state. I order to establish a relation between the plastic component of deformation and stress increment, an assumption is introduced, that the increment of plastic deformation is proportional to the stress gradient of size which is termed the plastic potential, so that:

$$\{d\varepsilon^p\} = \lambda \frac{\partial Q}{\{\partial \sigma\}} \tag{4}$$

In the term (4) Q is the function of the plastic potential and λ is a scalar which determines the value of plastic deformation increment. Considering that the gradient of the function is normal to this function, term (4) represents a normality principle in the theory of plasticity and can be represented by Fig. 5.

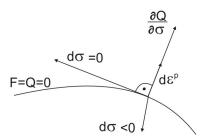


Fig. 5 Normality Principle in the plasticity theory

Term (4) represents relation between the increment of plastic deformation and stress in the plastic area, that is, it determines the direction of increment vectors of plastic deformations and is called the plastic yield law. If the surface of plastic potential Q and yield surface F coincide, as in Fig. 4, it is an associate plasticity so the increment vector of plastic deformation is normal to the yield curve.

In Fig. 4 it is visible that if the stress point is on the yield surface the stress increment $d\sigma$ can have any direction. If it is directed towards the yield surface F and neutral, that is, if it is tangent to eh yield surface, the unloading will have elastic deformations, while if it is directed outside the yield surface, the surface will change and elastic and plastic deformations will occur.

3. CLASSIC PLASTIC MODELS OF MATERIALS

In historical perspective, many idealized models of material behavior have been proposed, and even nowadays this is the field of much effort and progress, thus the number of models is extremely high, especially when those are models of the soil. Determination of the yield surface and failure stress envelope in the main stress space is huge and very

complex issue. For this reason, idealizations were introduced, so that the problem could be solved more easily and obtained solution applicable in practice. Many of the proposed models did not match the experimental results, so they are only interesting from the historical point of view. Here, only some of the models which were affirmed in practice will be presented.

The simple models which include the plastic properties in the description of material behavior are ideally elasto-plastic models. Working hardening or weakening is not present in them so all the yield surfaces in the spatial stress state are blended into one – final surface called the failure surface. Inside the failure surface, material behaves as linear-elastic, and when the stress state is on the yield surface a plastic deformation occurs whose direction and value are determined on the basis of potential function and failure conditions.

Ideal elasto-plastic models described by the classical failure theories approximate the curved failure envelope to the solids whose generatrices are straight lines. In the Fig. 6 some of the most famous classic ideally elasto-plastic models are presented. Surfaces of yield, that is, failure are defined by the criteria of: Tresca, Von Mises, Mohr-Coulomb and Drucker-Prager.

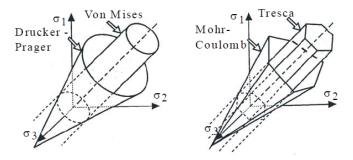


Fig. 6 Classic failure theories – spatial geometrical interpretation. [4]

Advantage of these models is their simplicity, easy determination of yield parameters and relatively simple application during software writing. Its shortcomings could be the unrealistically high dilatation and inability to imitate hysterestesis.

3.1. Tresca's yield criterion (1864)

On the basis of numerous experiments Tresca concluded that yield begins when maximum shear stress in some point reaches critical value. For the case when for the main stresses σ_1 , σ_2 , σ_3 is valid $\sigma_1 \ge \sigma_2 \ge \sigma_3$ this criterion can, according to [1] be expressed in a following way:

$$\tau_{\text{max}} = \frac{1}{2} (\sigma_1 - \sigma_3) = k_T \tag{5}$$

Where k_T is a yield constant experimentally determined depending on the type of material.

The value of the constant k_T can be determined from the test of uniaxial stress in dependence of the tensile yield stress σ_T . For this case of stress, $\sigma_1 = \sigma_T$, a $\sigma_1 = \sigma_2 = 0$, so that the following results from the term (5):

$$k_T = \frac{\sigma_T}{2} \tag{6}$$

The value of constant k_T can be determined from the pure shear test depending on the shear yield stress. In such case of stress $\sigma_1 = -\sigma_3 = \sigma_T$ a $\sigma_3 = 0$ so from the expression (5) it is:

$$k_T = \tau_T \tag{7}$$

Tresca's yield surface can, in the main stress space be represented by infinitely long hexagonal prism as in Fig. 7, that is by the term:

$$\max(|\sigma_1 - \sigma_2|, |\sigma_2 - \sigma_3|, |\sigma_3 - \sigma_1|) = k_T$$
(8)

The axis of the prism coincides with the spatial diagonal (hydrostatic axis) determined by the points $\sigma_1 = \sigma_2 = \sigma_3$.

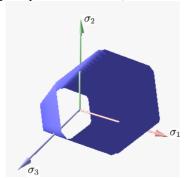


Fig. 7 Tresca's yield surface in the main stresses space

Material stays in the elastic area for as long as all three main stresses are approximately equal, regardless of how pressed or tensioned it is. This means that hydrostatic stress does not influence the yield, so all the intersections with octahedron planes (planes perpendicular to hydrostatic axis) are equal. An intersection with one octahedron plane is presented in Fig. 10. However, if the material is exposed to shearing, that is, if one of the main stresses becomes higher or lower than the rest, the stress will proceed to the yield surface and plastic deformations start to occur.

In the two dimensional space of main stresses σ_1 , σ_2 the Tresca's yield surface is represented by the Fig. 8.b.

3.2. Von Mises' yield criterion (1913)

Von Mises concludes that material starts to yield when the second invariant of stress deviator reaches the critical value, that is according to [1]:

$$I_2 = -k_M^2 \tag{9}$$

The term (9) can be written in its derived form, in the following way:

$$I_2 = -\frac{1}{6} [(\sigma_{11} - \sigma_{22})^2 + (\sigma_{22} - \sigma_{33})^2 + (\sigma_{33} - \sigma_{11})^2 + 6(\sigma_{12}^2 + \sigma_{23}^2 + \sigma_{31}^2)] = -k_M^2$$
 (10)

The value of the yield constant k_M can be determined from the test of uniaxial stress where $\sigma_{11} = \sigma_T$ while other stress components are equal to zero. Then the term (9) is transformed into:

$$\frac{1}{3}\sigma_T^2 = k_M^2 \text{ that is } k_M = \frac{\sigma_T}{3}$$
 (11)

On the basis of the previous statements, the yield surface is:

$$F = -\frac{\sigma_T}{3} \tag{12}$$

For this model too, the value k_M can be determined form the pure shear test where $\sigma_{11} = \sigma_{22} = \sigma_{33} = \sigma_{23} = \sigma_{31} = 0$, so that the result is:

$$k_M = \tau_T \tag{13}$$

In the three-dimensional space of main stresses σ_1 , σ_2 , σ_3 von Mises' yield surface is circular cylinder whose axis, as in the Tresca's yield surface, coincides with hydrostatic axis. Intersection of both yield surfaces with the planes σ_1 , σ_2 is presented in the Fig. 8.b.

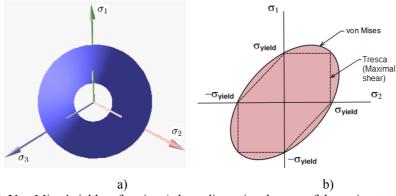


Fig. 8 Von Mises' yield surface in: a) three-dimensional space of the main stresses σ_1 , σ_2 , σ_3 and b) in two-dimensional space of the main stresses σ_1 , σ_2

The tests on many metals demonstrated that Von Mises' criterion is in much better accordance with the results of experiments in respect to the Tresca's criterion. In spite of that, Tresca's criterion is much more frequently applied, primarily due to the simplicity of application.

3.3. Mohr-Coulomb yield criterion (1882)

The starting point here is long known Coulomb condition of soil failure (1773) which states that the failure occurs when the shear stress in a plane reaches the value:

$$\tau = c + \sigma \cdot tg\phi \tag{14}$$

In the term (14) c and φ are cohesion and angle of interior soil friction, where the adopted convention is that the pressure stress is positive.

Mohr (1882) determined that equation (14) graphically represents a tangent on the larges circle of main stresses. The tangent is simultaneously the failure stress envelope, which can be seen in Fig. 9. Through the main stresses σ_1 and σ_3 it is expressed, according to [3], in the following way:

$$(\sigma_1 - \sigma_3) = (\sigma_1 + \sigma_3)\sin\varphi + 2c \cdot \cos\varphi \tag{15}$$

In the three-dimensional space of the main stresses σ_1 , σ_2 , σ_3 the Mohr-Coulomb failure criterion is usually expressed in the following way:

$$\pm \frac{\sigma_1 - \sigma_2}{2} = \frac{\sigma_1 + \sigma_2}{2} \sin \varphi + c \cdot \cos \varphi$$

$$\pm \frac{\sigma_2 - \sigma_3}{2} = \frac{\sigma_2 + \sigma_3}{2} \sin \varphi + c \cdot \cos \varphi$$

$$\pm \frac{\sigma_3 - \sigma_1}{2} = \frac{\sigma_3 + \sigma_1}{2} \sin \varphi + c \cdot \cos \varphi$$
(16)

In the main stress space, the yield surface is irregular pyramidal surface with six faces, represented in Fig. 6, and its intersection with the octahedron plane is presented in the Fig. 10. This ideally plastic model has been often used for presentation of cohesive-granular materials such as soil, rock and concrete. When $\sigma_1 = \sigma_2 = \sigma_3$ from the expression (16) results the "apex" of the pyramid is on the spatial diagonal in the point $\sigma_1 = \sigma_2 = \sigma_3 = c \cdot ctg\varphi$. When the materials in question are without cohesion, that is, when c = 0, the "apex" of the pyramid is displaced to the origin of the coordinate system.

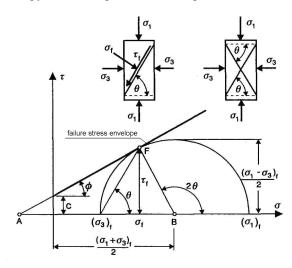


Fig. 9 Mohr's failure stress diagram [3]

The shape of the pyramid with six faces in the octahedron plane is defined by eh angle of internal friction of material which can be between 0^0 and 90^0 . When $\varphi = 0^0$ Mohr-Coulomb model is transformed into the Tresca's model and in the case $\varphi = 90^0$ it becomes the Rankine's model with triangular deviator cross section, which is illustrated in the Fig. 10.

Mohr-Coulomb's model successfully reproduces the trait that shear strength depends on the level of mean stress. On the edges of the yield pyramid there are singularities, that is, the direction of plastic deformation increment vector is not unequivocally determined. The problem of change of yield surface shape is a frequent problem encountered by the model in practice, and occurs also in the Tresca's yield surface. It is usually solved by rounding the pyramid angles.

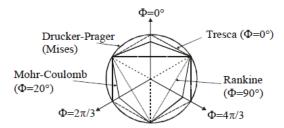


Fig. 10 Intersections of yield surfaces with octahedron (deviator) plane [4]

3.4. Drucker-Prager's yield criterion (1952)

Like Mohr-Coulomb model, it belongs to elasto-plastic models where on the yield surface there are only plastic deformations. It was made as a simplification of Mohr-Coulomb and expansion of von Mises models. In von Mises' expression, the influence of hydrostatic stress component on yield is expressed by introduction of an additional member:

$$\alpha I_1 + I_2^{1/2} = k \tag{17}$$

In the main stress space, the yield surface is a cone around the hydrostatic axis, and in the octahedron surface it is a circle with \sqrt{k} diameter. In case when Drucker-Prager circle passes through the external apexes of Mohr-Coulomb hexagon, the parameters α and k have the values:

$$\alpha = \frac{2\sin\varphi}{\sqrt{3}(3-\sin\varphi)} \text{ and } k = \frac{6c\cdot\cos\varphi}{\sqrt{3}(3-\sin\varphi)}$$
 (18)

When it touches on the interior of Mohr-Coulomb surface of yield, then they have the following values:

$$\alpha = \frac{2\sin\varphi}{\sqrt{3}(3+\sin\varphi)} \text{ and } k = \frac{6c\cdot\cos\varphi}{\sqrt{3}(3+\sin\varphi)}$$
 (19)

5. CONCLUSION

Establishment of constitutive relations in order to define the way of behavior of materials under external load is a complex process even in the simplest ideally elastic and ideally plastic models. Considering that they, nowadays cannot describe in a satisfactory way the real behavior of materials in all conditions of usage, the model theory developed in the direction of introduction of more quality elasto-plastic models into calculation. Further development is a continuous process, where new models are created and which describe the nature of the materials progressively better. Apart from that the development of computer technology made possible that many of them are practically applied in calculations. In the following paper, which is a continuation of this one, the models used in the Ansys software package.

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MATEMATIČKO MODELIRANJE MATERIJALNO NELINEARNIH PROBLEMA U ANALIZI KONSTRUKCIJA (I DEO – TEORIJSKE OSNOVE)

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Modeli materijala opisuju način njihovog ponašanja pri opterećivanju. U radu je predstavljen razvoj modela počevši od najjednostavnijih linearno elastičnih i krutoplastičnih do znatno složenijih idealno elastoplastičnih. Definisani su osnovni pojmovi teorije plastičnosti kao što su kriterijum i površ tečenja (popuštanja), zakon ojačanja, zakon plastičnog tečenja i uslov normalnosti. Posebno su dati kriterijumi tečenja Tresca-e, Von Mises-a, Mohr-Coulomb-a i Drucker-Prager-a.

Ključne reči: materijalni modeli, plastičnost, kriterijum tečenja, površ tečenja